

# HII-CHI-mistry: A collection of python scripts for the analysis of emission lines in ionized gaseous nebulae

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## 1. General description

### 1.1. What is HII-CHI-MISTRY and where can it be found?

Ionized gaseous nebulae are ubiquitous objects in the Universe that provide valuable information about the galaxies where they reside through their bright emission lines. Among the various properties that can be derived from these nebulae are the relative chemical abundances of observed ions, the excitation of the gas, and the effective hardness of the ionizing source.

HII-CHI-MISTRY (hereafter, HCM) is a collection of Python scripts designed to analyze observational data from different bright emission lines measured in the ultraviolet, optical, or infrared ranges of the spectrum. By comparing these observations with predictions from extensive grids of photoionization models, HCM provides estimates and associated uncertainties for several derived properties, adapting the accuracy of the solutions based on the input lines and their uncertainties.

HCM offers three main advantages for the calculation of chemical abundances compared to other model-based solutions:

1. The results are fully consistent with the direct method, which relies on the prior determination of the electron temperature. This holds true even when the flux of an auroral line, such as  $[\text{O III}] \lambda 4363 \text{ \AA}$ , is not used (see Pérez-Montero, 2014). 2. HCM provides consistent solutions regardless of the set of input emission lines and their errors. This is particularly useful when comparing results for different sets of objects observed at different spectral ranges or redshifts. 3. HCM provides in a first step solutions for N/O (in the optical and infrared) or C/O (in the ultraviolet), allowing a later more accurate determination of O/H using N or C without assuming any prior relationship between O/H and N/O or C/O.

In this tutorial, I explain the features of the different versions of the code, their usage instructions, limitations, and expected future improvements. There are five versions of HCM depending on the spectral range and whether they are used for determining chemical abundances or calculating the hardness of the ionizing source. These versions are:

- HII-CHI-MISTRY: The original HCM package described in Pérez-Montero, 2014 for analyzing gaseous nebulae ionized by different sources of ionization including massive stars, active galactic nuclei, post-AGB stars or advected-dominated flows. It can be used to derive the total oxygen abundance ( $12 + \log(\text{O/H})$  (hereafter O/H), the total sulphur abundance ( $12 + \log(\text{S/H})$ , hereafter S/H), the nitrogen-to-oxygen chemical abundance ratio ( $\log(\text{N/O})$ , hereafter N/O), and the ionization parameter ( $\log U$ ) using optical emission lines from  $[\text{O II}] \lambda 3727 \text{ \AA}$  up to  $[\text{S III}] \lambda 9069,9532 \text{ \AA}$ .
- HII-CHI-MISTRY-UV: Described in Pérez-Montero and Amorín, 2017 for star-forming objects, and in Pérez-Montero et al., 2023a for the NLR of AGN. It calculates O/H, the carbon-to-oxygen chemical abundance ratio ( $\log(\text{C/O})$ , hereafter C/O), and  $\log U$ , using ultraviolet emission lines from  $\text{Ly}\alpha \lambda 1216 \text{ \AA}$  up to  $\text{C III} \lambda 1909 \text{ \AA}$ .

- HII-CHI-MISTRY-IR: Described in [Fernández-Ontiveros et al., 2021](#) for star-forming objects, and in [Pérez-Díaz et al., 2022](#) for the NLR of AGNs. It estimates O/H, S/H, N/O, and  $\log U$  using infrared emission lines from  $\text{Br}\alpha \lambda 4.05 \mu\text{m}$  up to  $[\text{N II}] \lambda 205 \mu\text{m}$ .
- HII-CHI-MISTRY-TEFF: It calculates the equivalent effective temperature ( $T_{\text{eff}}$ ) and  $\log U$  using as input O/H and ratios of consecutive optical and UV emission lines, as described in [Pérez-Montero et al., 2019b](#). It can also be used to derive the fraction of absorbed ionizing photons for objects with  $\text{He II} \lambda 4686 \text{\AA}$ , as described in [Pérez-Montero et al., 2020](#).
- HCM-TEFF-IR: It estimates  $T_{\text{eff}}$  and  $U$  based on near- and mid-infrared emission lines, as described in [Pérez-Montero et al., 2024](#). It can also calculate the parameter  $\alpha_{OX}$ , indicative of the hardness of the incident SED of AGN taking advantage of the very high-excitation lines present in this regime (see [Pérez-Montero et al., 2025](#)..

HCm has been originally written in python v. 2.7, but from version 5 is only compatible with python 3. It can be downloaded from the HCm webpage at <http://home.iaa.es/~epm/HII-CHI-mistry.html>. In this document, I describe features for the public versions of HCm 5.5 and equivalent versions, which require the Python library `numpy`. Each package contains the corresponding script (with a `.py` extension), an `ASCII` file with instructions, an example input text file, and different libraries containing model information and template files for constrained sub-grids used by the code in the absence of certain emission lines. All files are included in a compressed `tgz` file, which can be uncompressed using the following command in a terminal prompt:

```
> tar xvfz HCm_v5.5.tar.gz
```

## 1.2. How does the code work?

All versions of HCm employ a Bayesian-like approach to derive chemical abundances, the ionization parameter, or  $T_{\text{eff}}$ . In brief, for a given property  $X$ , the final result is calculated as follows:

$$X_f = \frac{\sum_i X_i / \chi_i}{\sum_i 1 / \chi_i} \quad (1.1)$$

where  $X_f$  is the final result,  $X_i$  are the input values from each model in the grid, and  $\chi_i$  are the weights assigned to each model, calculated as the square difference of the observed and predicted values for specific emission-line ratios:

$$\chi_i^2 = \sum_j \frac{(O_j - T_{ji})^2}{O_j} \quad (1.2)$$

where  $O_j$  and  $T_{ji}$  are the observed and model-based values, respectively, for the considered emission-line dependent ratios. These ratios are described in each of the papers on the different versions of HCm as a function of the input emission lines. Notice that, contrary to a pure bayesian method, in which the model with the lowest  $\chi^2$  is selected, this method provides a more realistic estimation of the derived properties, even if the nominal values do not exactly match any of the models, but giving estimates for the corresponding uncertainties.

The errors assigned to each result are calculated as the square sum of two sources. The first is the dispersion of all results obtained through Monte Carlo iterations made

by randomly perturbing the input nominal values with the input observational errors. The second is the mean of all intrinsic uncertainties assigned to the Bayesian process, calculated as:

$$\Delta X^2 = \frac{\sum_i (X_f - X_i)^2 / \chi_i}{\sum_i 1 / \chi_i} \quad (1.3)$$

For the versions of HCm aimed at calculating chemical abundances, the first iteration throughout the entire chosen grid of models provides an estimation for the abundance of a secondary ion (N for HCm and HCm-IR, or C for HCm-UV), relative to oxygen. These estimations are based on emission-line ratios that are not very sensitive to  $\log U$ . This allows the grid to be constrained in a second iteration, once N/O or C/O are fixed, to calculate both O/H and  $\log U$  using N or C lines without any prior assumption about the relationship between O and the respective secondary element.

In the case of HCm-TEFF, both for the versions for the optical and the IR, this previous iteration is not performed. Instead the grid of models is interpolated to fix the value of O/H to the input value in order to minimize the dependence of both  $T_{\text{eff}}$  and  $U$  on metallicity.

## 2. HCm in the optical

### 2.1. Running the program

To run the HCm code version 6.0, type the following command in the terminal prompt:

```
> python HCm_v6.0.py
```

A description of the code will appear on-screen, prompting for the input file containing observational information. Alternatively, the input file name and the desired number of iterations for the Monte Carlo simulation can be specified when invoking the code:

```
> python HCm_v6.0.py input.txt 100
```

If the number of Monte Carlo iterations is not specified, it defaults to 25.

### 2.2. The input file

The input file must be an ASCII file, with as many rows as the number of objects or pointings for which the calculations are to be performed. Each column represents the identification for each row and the reddening-corrected fluxes relative to H $\beta$  with their corresponding errors. In previous versions of HCm, all columns needed to be introduced, but from version 5.0 onwards, a first row with the labels of the introduced columns is required, and the order is not essential. The file can also include other columns not identified by the code. The columns dedicated to the observational relative errors are optional. The labels for the emission lines in this version are:

- ID: To identify each row with a name.
- OII\_3727 and eOII\_3727: For  $[[\text{O II}]] \lambda 3727 \text{ \AA}$  and its error. Assumed to be the sum of  $\lambda 3726 \text{ \AA}$  and  $\lambda 3729 \text{ \AA}$  in resolved spectra.
- NeIII\_3868 and eNeIII\_3868: For  $[[\text{Ne III}]] \lambda 3868 \text{ \AA}$  and its error.
- OIII\_4363 and eOIII\_4363: For  $[[\text{O III}]] \lambda 4363 \text{ \AA}$  and its error.
- OIII\_4959 and eOIII\_4959: For  $[[\text{O III}]] \lambda 4959 \text{ \AA}$  and its error.

- OIII\_5007 and eOIII\_5007: For  $[[\text{O III}]] \lambda 5007 \text{ \AA}$  and its error. If only one of the two strong  $[[\text{O III}]]$  lines is provided, the code assumes the theoretical ratio between them.
- NII\_5755 and eNII\_5755: For  $[[\text{N II}]] \lambda 5755 \text{ \AA}$  and its error.
- SIII\_6312 and eSIII\_6312: For  $[[\text{S III}]] \lambda 6312 \text{ \AA}$  and its error.
- NII\_6584 and eNII\_6584: For  $[[\text{N II}]] \lambda 6584 \text{ \AA}$  and its error.
- SII\_6716 and eSII\_6716: For  $[[\text{S II}]] \lambda 6716 \text{ \AA}$  and its error.
- SII\_6731 and eSII\_6731: For  $[[\text{S II}]] \lambda 6731 \text{ \AA}$  and its error. Note that the code will only use the  $[[\text{S II}]]$  lines if both of them or their sum is provided.
  - Alternatively, the sum of these two lines can be given as SII\_6725 and eSII\_6725.
  - OII\_7325 and eOII\_7325: For  $[[\text{O II}]] \lambda\lambda 7319+7330 \text{ \AA}$  and its error.
  - SIII\_9069 and eSIII\_9069: For  $[[\text{S III}]] \lambda 9069 \text{ \AA}$  and its error.
  - SIII\_9532 and eSIII\_9532: For  $[[\text{S III}]] \lambda 9532 \text{ \AA}$  and its error. If only one of the two strong  $[[\text{S III}]]$  lines is provided, the code assumes the theoretical ratio between them.

### 2.3. Selecting the grid of models

It is not necessary to write all columns as not all the lines are mandatory and the code can work to provide a solution using only a set of them. If the input file is correct, the code will prompt you to select the grid of models to perform the calculation:

- (1) POPSTAR with Chabrier IMF, age = 1 Myr
- (2) BPASS v.2.1 a\_IMF = 1.35, Mup = 300, age = 1 Myr with binaries
- (3) AGN, double component,  $\alpha(\text{UV}) = -1.0$
- (4) pAGB NLTE models (Rauch 2003)
- (5) Advection-dominated Accretion Flow for AGN (ADAF, Nemmen et al. 2014)

Other SED

- (6) Different library

Choose SED of the models:

All grids have been calculated using the code CLOUDY v.17 (Ferland et al., 2017) assuming a central ionizing source and a plane-parallel geometry. Grid 1 is described in Pérez-Montero, 2014 and is calculated using POPSTAR (Mollá et al., 2009) cluster model atmospheres with an instantaneous burst at an age of 1 Myr, assuming a Chabrier, 2003 initial mass function (IMF) and a constant electron density of  $100 \text{ cm}^{-3}$ . Grid 2 is described in Pérez-Montero et al., 2021 and was calculated with cluster model atmospheres from BPASS v.2.1 (Eldridge et al., 2017), assuming an instantaneous burst at an age of 1 Myr, an IMF with slope  $x = -1.35$ , an upper mass limit of  $300 \text{ M}_\odot$ , and binaries. The gas in these models has an electron density of  $100 \text{ cm}^{-3}$ . This grid is recommended for Extreme Emission Line Galaxies (EELGs).

Option 3 can be used to derive chemical abundances in the Narrow Line Regions (NLR) of AGN and includes several grids described in Pérez-Montero et al., 2019a. These assume a double-peaked power law spectral energy distribution with a parameter  $\alpha_{UV} = -1.0$ . The assumed electron density is  $500 \text{ cm}^{-3}$ . The code will ask for the value of  $\alpha_{OX}$  with possible values from -0.8 to -2.0 in steps of 0.2, and the criterion used to stop the models (at an outer radius when the fraction of free electrons is 98% or 2%).

Models for post-AGB stars were calculated using libraries from Rauch, 2003 covering values for  $T_{\text{eff}} = 5, 10, \text{ and } 15 \cdot 10^4 \text{ K}$ , while models simulating advected-dominated flows (ADAFs) were taken from Nemmen et al., 2014. These two sets of models are available only from version 5.5 and they are described in Pérez-Díaz et al., 2025.

All default grids cover an input O/H value in the range [6.9, 9.1] in bins of 0.1 dex and N/O in the range [-2.0, 0.0] in bins of 0.125 dex. Log  $U$  is covered with a resolution

of 0.25 dex in the range [-4.0, -1.5] for grids for star-forming regions (i.e., 1 and 2) and in the range [-2.5, -0.5] for grids for AGN. From version 6.0, star-forming models have been also calculated assuming a variable S/O ratio ranging from -0.6 to +0.6 dex around the solar ratio in bins of 0.2 dex. These libraries are used by the code to provide an estimation of S/H (see Pérez-Montero et al, submitted).

The user may also use his/her model libraries by including them as a text file in the correct format in the folder `Libraries_opt`. If a different library is introduced by the user (4), the code will automatically check if the file has the correct format. If some information is missing, the code will warn the user about the missing columns.

Once the grid of models has been selected, the code will ask whether to use interpolation between the models in the grid:

**Choose models [0] No interpolation [1] Interpolation:**

In the interpolated mode, the code makes a linear interpolation of all variables, increasing the resolution by a factor of 10. This prevents the results from clustering around certain grid points. However, this mode slows the calculation time.

As a final step, the program will ask for the constraint laws to limit the grids. This is necessary when a limited set of emission lines is provided and the models must make assumptions about the relations between O/H and N/O (when N/O cannot be derived) or O/H and  $U$  (when no auroral line is provided, so excitation is used to estimate metallicity). For AGNs, there is a degeneracy of the O2O3 emission-line ratio, as discussed in Pérez-Montero et al., 2019a. To derive  $\log U$ , the code will also ask for the preferred  $\log U$  range for the calculations between high-ionization ( $\log U \geq -2.5$ ) or low-ionization ( $\log U \leq -2.5$ ).

**Select a file with the constraint laws to limit the grid of models when the measurement of a quantity is impossible without any relation.**

#### Default constraints

- (1) Constraints for Star-Forming Galaxies
- (2) Constraints for Extreme Emission Line Galaxies
- (3) Constraints for AGNs (no restriction on the ionization parameter)
- (4) Constraints for high-ionization AGNs ( $\log(U) > -2.5$ )
- (5) Constraints for low-ionization AGNs ( $\log(U) < -2.5$ )

#### Other constraints

- (6) Different constraint file

The user can also use his/her own constraint assumptions in a text file with the correct format in the folder `Constraints`. If option 6 is chosen, the code will check the file for the correct format. If some information is missing, the code will warn the user about the missing columns.

After this process, the program will summarize the different grids that will be used and the number of models for each grid. From version 5.4 onwards, this process can be sped up by changing the variable `interactive` from `True` to `False` in the script and editing the model parameters and constraints directly in the code.

#### 2.4. Results

The code will display the results of the calculations for each row on the screen, along with the task's completeness ratio. At the end, an additional ASCII file containing all results

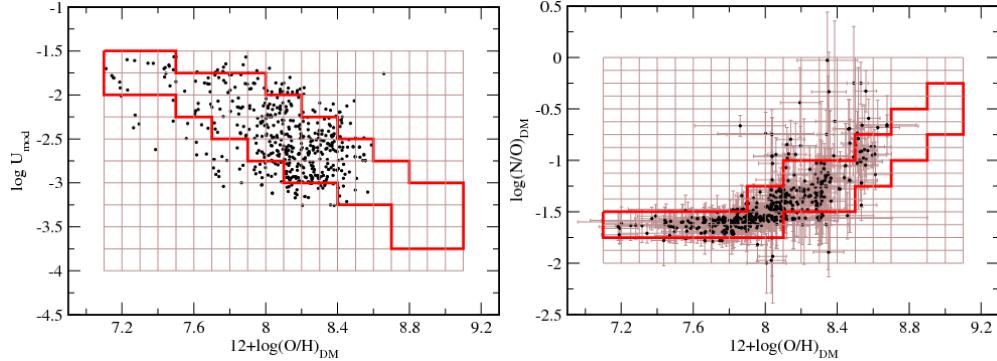


FIGURE 1. Left: relation between total oxygen abundance and ionization parameter for the sample studied in Pérez-Montero, 2014. The solid red line encompasses the most probable combination of parameters occupied by the objects. Right: empirical relation between  $12+\log(\text{O/H})$  and  $\log(\text{N/O})$  with the region occupied by the grid of models in the absence of observational information to constrain N/O.

will be generated, named the same as the input file but with the suffix `.hcm-output.dat`. The first column lists the identification of each row. If this was not specified in the input file, a number will be assigned. The next columns of this file contain the emission-line fluxes used as input, with their corresponding errors.

The next column is an index that indicates whether the whole grid of models has been used or if a constrained grid has been used instead. This depends mainly on the introduced emission lines for each row, as a limited observational set implies additional assumptions to calculate ionic abundances. However, as described above, these constraints can be chosen based on the models, or the user can use their own. An index of 1 indicates that the whole grid of models has been used. This is done when an estimation of the electron temperature, similar to the direct method (see for instance Pérez-Montero, 2017), can be made because both auroral and nebular emission lines are present (e.g.,  $[\text{O III}]$  4363 Å with 4959, 5007 Å,  $[\text{N II}]$  5755 Å with 6584 Å,  $[\text{O II}]$  7325 Å with 3727 Å, or  $[\text{S III}]$  6312 Å with 9069, 9532 Å). If any auroral line is not given, which is common in faint or metal-rich nebulae (Pérez-Montero and Díaz, 2005), additional assumptions should be made.

Index 2 represents the case when the electron temperature cannot be estimated, and the code assumes an empirical law between O/H and  $\log U$  (see Figure 1). This relation assumes that metal-poor objects generally exhibit higher excitation, while metal-rich objects have lower excitation. This is the same assumption behind the use of many high-to low-excitation emission line flux ratios (e.g., O3O2, O3N2) to derive chemical abundances. Finally, index 3 denotes the use of a grid constrained by assuming an empirical relation between O/H and N/O, as shown in Figure 1. This is necessary when N/O cannot be calculated independently in a first iteration using emission-line ratios such as N2O2, O3N2, or N2S2. This assumption is behind all strong-line calibrations based on  $[\text{N II}]$  and implies a constant N/O value for low metallicities due to primarily primary production of N, and an increasing N/O with O/H when secondary production of N is assumed. This assumption can lead to non-negligible deviations from the real O/H if N/O does not lie in the expected regime, as discussed in Pérez-Montero and Contini, 2009.

The final 8 columns of the output file give the results for O/H, S/H, N/O, and  $\log U$  with their corresponding errors. If no solution is found, O/H and  $\log U$  are denoted by

TABLE 1. Mean offsets and standard deviation of the residuals of the resulting properties derived by HCm when the model from POPSTAR emission lines are used as input as a function of the used emission-line ratios.  $[[\text{O III}]]_a$  stands for the  $[[\text{O III}]]$  auroral line at  $\lambda 4363 \text{ \AA}$  and  $[[\text{O III}]]_n$  for the nebular lines at  $\lambda 4959, 5007 \text{ \AA}$ .

Used lines	Grid	O/H		N/O		$\log U$				
		Mean	$\Delta$	$\sigma$ res.	Mean	$\Delta$	$\sigma$ res.	Mean	$\Delta$	$\sigma$ res.
All lines	1	+0.03	0.09	-0.07	0.08	+0.06	0.15			
$[[\text{O III}]]_a, [[\text{O III}]]_n, [[\text{N II}]], [[\text{S II}]]$	1	+0.02	0.09	-0.01	0.07	+0.06	0.15			
$[[\text{O II}]], [[\text{O III}]]_n, [[\text{N II}]], [[\text{S II}]]$	2	-0.04	0.08	+0.00	0.03	+0.05	0.10			
$[[\text{O III}]]_n, [[\text{N II}]], [[\text{S II}]]$	2	-0.03	0.07	+0.00	0.04	+0.05	0.10			
$[[\text{O II}]], [[\text{O III}]]_n, [[\text{N II}]]$	2	-0.04	0.24	+0.00	0.23	+0.03	0.09			
$[[\text{N II}]], [[\text{S II}]]$	2	+0.00	0.19	+0.00	0.04	+0.02	0.20			
$[[\text{O III}]]_n, [[\text{N II}]]$	3	-0.04	0.08	—	—	+0.07	0.13			
$[[\text{N II}]]$	3	+0.01	0.15	—	—	+0.03	0.21			
$[[\text{O II}]], [[\text{O III}]]_n$	3	+0.00	0.01	—	—	+0.00	0.02			
$[[\text{O II}]], [[\text{Ne III}]]$	3	-0.01	0.02	—	—	-0.01	0.03			

0, and N/O is denoted by -10. If the code encounters an error for a specific row of the input, '9999' will be inserted in the output. In this case, the user should review the input lines or contact us for possible inconsistencies.

The results and consistency of the abundances derived from HCm for star-forming objects, whose abundances were derived from the direct method (Pérez-Montero, 2014), and for NLR in AGN as derived using detailed photoionization models (Pérez-Montero et al., 2019a), are well discussed. Nevertheless, in Table 1, I provide a list of the mean offsets and the standard deviation of the residuals of the resulting O/H, N/O, and  $\log U$  as compared with the input values from the models as a function of the used emission lines. These values illustrate how the code can recover the abundances from the values using only the emission lines as input, but cannot be taken as true uncertainties.

### 3. HCm in the ultraviolet

The HII-CHI-MISTRY-UV code (hereafter HCm-UV) is similar to the optical version but uses a different set of emission lines in the ultraviolet regime and estimates C/O instead of N/O in the first iteration. It also allows including some emission lines in the optical range to estimate the abundance based on an emission-line ratio sensitive to the electron temperature (e.g.,  $[[\text{O III}]] \lambda 5007/\lambda 1665$ ), although this may introduce larger inaccuracies due to reddening uncertainties. This code is described in Pérez-Montero and Amorín, 2017 for star-forming objects and in Pérez-Montero et al., 2023a for the NLR in AGN.

#### 3.1. Running the Program and Preparing the Input File

Similar to the optical version, it can be executed from the terminal using Python, allowing the script to be run directly or with a specified input file and iteration count for Monte Carlo simulations:

```
> python HCm-UV_v5.0.py input.txt 100
```

The input file is an `ascii` file where the first row must contain labels for each column, including:

- ID: Identification name for each row.
- Ly $\alpha$ \_1216 and eLy $\alpha$ \_1216: Ly $\alpha$   $\lambda$  1216 Å and its error.
- NV\_1239 and eNV\_1239: N v]  $\lambda$  1239 Å and its error.
- CIV\_1549 and eCIV\_1549: CIV]  $\lambda$  1549 Å and its error.
- HeII\_1640 and eHeII\_1640: He II $\lambda$  1640 Å and its error.
- OIII\_1665 and eOIII\_1665: [O III]  $\lambda$  1665 Å and its error. This includes all lines of the [O III] multiplet between 1660 Å and 1666 Å.
- CIII\_1909 and eCIII\_1909: C III] 1909 Å and its error. Again, this includes all lines of the C III] multiplet.
  - OII\_3727 and eOII\_3727: [[O II]] 3726+3729 Å and its error.
  - NeIII\_3869 and eNeIII\_3869: [[Ne III]] 3869 Å and its error.
  - OIII\_4363 and eOIII\_4363: [[O III]] 4363 Å and its error.
  - Hb\_4861 and eHb\_4861: H $\beta$   $\lambda$  4861 Å and its error.
  - OIII\_4959 and eOIII\_4959: [[O III]] 4959 Å and its error.
  - OIII\_5007 and eOIII\_5007: [[O III]]  $\lambda$  5007 Å and its error.

Subsequent rows should include identification names and extinction-corrected fluxes in arbitrary units. Notice that, contrary to the version for the optical, the H $\beta$  flux must be introduced separately to let some observables, such as R23, to be considered. A value of zero can be used for missing values, and not all columns are required to estimate O/H, C/O, or  $U$ .

If the input file is correct, the code will prompt for the grid of models and the use of interpolation. Version 5.0 offers models from POPSTAR, BPASS, and double composite AGN. For AGN, the code will ask for the assumed  $\alpha_{OX}$  (from -0.8 to -2.0 in bins of 0.2), the stopping criterion for the models (2% or 98

In the next step, the code will prompt for model interpolation to improve resolution and the choice of templates considering the assumed relation between O/H,  $\log U$ , and C/O in case that some of the input emission lines necessary to derive these quantities are missing. These templates can be edited by the user in the appropriate folder. As in the optical version, interactivity can be disabled by setting the variable `interactive` to `False`.

### 3.2. Output File and Analysis of the Results

Once the grid of models is selected, the code will display the results for each input row on the screen, along with the completeness ratio. At the end, it generates an output file in `ascii` format containing the results, including all information used for the calculations. The first column is the corresponding ID, and the next columns will be the input emission lines along with their errors.

The following column is an index indicating whether the complete model grid was used (index 1), which occurs only when at least two [O III] lines among  $\lambda$  1665, 4363 and 5007 Å are provided to estimate the electron temperature. Index 2 denotes a grid with an implicit relation between O/H and  $\log U$  as shown in Figure 1. Finally, index 3 denotes a grid constrained by assuming a relation between C/O and O/H, similar to that shown in Figure 1 for N/O and O/H, considering a fixed C/N to the solar ratio. This index is applied when C/O cannot be calculated using the C3O3 or C3Ne3 emission line ratios, depending on C III] 1909 Å and [O III]  $\lambda$  1665 Å or [[Ne III]] 3869 Å, respectively, but an estimate for both O/H and  $\log U$  can be given using C lines. In any case, as commented above, this relation can be changed by the user in the folder `Libraries_uv`. The code, for instance, also provides the relation derived for EELGs given by Pérez-Montero et al., 2021.

The final six columns give the results for O/H, C/O, and  $\log U$  with their corresponding

TABLE 2. Mean offsets and standard deviation of the residuals of the resulting properties derived by HCm-UV when the model from POPSTAR emission lines are used as input as a function of the used emission-line ratios.

Used lines	Grid	O/H			C/O			$\log U$		
		Mean	$\Delta$	$\sigma$ res.	Mean	$\Delta$	$\sigma$ res.	Mean	$\Delta$	$\sigma$ res.
All lines	1	+0.02	0.28	-0.05	0.08	+0.04	0.25			
$\text{Ly}\alpha, \text{C IV}, \text{He II}, [\text{O III}], \text{C III}$	2	+0.02	0.27	-0.06	0.19	+0.05	0.10			
$\text{C IV}, \text{He II}, [\text{O III}], \text{C III}$	2	+0.10	0.30	-0.14	0.13	-0.11	0.20			
$\text{Ly}\alpha, \text{C IV}, \text{He II}, \text{C III}$	3	+0.00	0.02	—	—	+0.00	0.01			
$\text{C IV}, \text{He II}, \text{C III}$	3	-0.01	0.03	—	—	+0.00	0.02			

errors. As in the case of the optical version, if no solution can be found for both O/H and  $\log U$ , these are denoted as 0 in the file, and -10 for C/O. A value of 9999 appears when an error is detected in a specific row.

Table 2 lists the mean offsets and the standard deviation of the residuals when we use as input for the code the same predictions from the models of the grid, when POPSTAR models are used as a function of the different combination of emission lines that lead to a solution. A similar table for the results for AGN can be found in Pérez-Montero et al (2023).

#### 4. HCm in the infrared

The program HII-CHI-MISTRY-IR (hereafter HCm-IR) calculates the total oxygen and sulphur abundance, nitrogen-to-oxygen ratio, and ionization parameter using a set of observed emission lines in the mid-infrared. It is described in Fernández-Ontiveros et al., 2021. Since version 3.0, the program also supports calculations for AGN narrow-line regions (NLR), as described in Pérez-Díaz et al., 2022 and Pérez-Díaz et al., 2024, along with other model SEDs, such as ADAF and post-AGB stars.

##### 4.1. Running the Program and Preparing the Input File

Similar to other versions, HCm-IR can be run from the terminal using Python, and the input text file and the number of Monte Carlo iterations can be specified in the same command:

```
> python HCm-IR_v3.3.py input.txt 50
```

The input file should be in `ascii` format, with the first row containing the row identifiers and emission line labels, using the following accepted labels:

- ID: Identification name for each row.
- HI\_4m and eHI\_4m: Pa $\alpha$  at  $\lambda$  4.07  $\mu\text{m}$  and its error.
- ArII\_7m and eArII\_7m: [[Ar III]] at  $\lambda$  6.98  $\mu\text{m}$  and its error.
- HI\_7m and eHI\_7m: Br $\alpha$  at  $\lambda$  7.46  $\mu\text{m}$  and its error.
- ArV\_8m and eArV\_8m: [[Ar IV]] at  $\lambda$  7.90  $\mu\text{m}$  and its error.
- ArIII\_9m and eArIII\_9m: [[Ar III]] at  $\lambda$  8.99  $\mu\text{m}$  and its error.
- SIV\_10m and eSIV\_10m: [[S IV]] at  $\lambda$  10.5  $\mu\text{m}$  and its error.
- HI\_12m and eHI\_12m: H $\alpha$  at  $\lambda$  12.46  $\mu\text{m}$  and its error.
- NeII\_12m and eNeII\_12m: [[Ne II]] at  $\lambda$  12.8  $\mu\text{m}$  and its error.
- ArV\_13m and eArV\_13m: [[Ar V]] at  $\lambda$  13.1  $\mu\text{m}$  and its error.

- `NeV_14m` and `eNeV_14m`:  $[\text{Nev}]$  at  $\lambda 14.9 \mu\text{m}$  and its error.
- `NeIII_15m` and `eNeIII_15m`:  $[[\text{Ne III}]]$  at  $\lambda 15.5 \mu\text{m}$  and its error.
- `SIII_18m` and `eSIII_18m`:  $[[\text{S III}]]$  at  $\lambda 18.8 \mu\text{m}$  and its error.
- `NeV_24m` and `eNeV_24m`:  $[\text{Nev}]$  at  $\lambda 24.3 \mu\text{m}$  and its error.
- `OIV_25m` and `eOIV_25m`:  $[\text{OIV}]$  at  $\lambda 25.9 \mu\text{m}$  and its error.
- `SIII_33m` and `eSIII_33m`:  $[[\text{S III}]]$  at  $\lambda 33.7 \mu\text{m}$  and its error.
- `OIII_52m` and `eOIII_52m`:  $[[\text{O III}]]$  at  $\lambda 52 \mu\text{m}$  and its error.
- `NII_57m` and `eNII_57m`:  $[[\text{N II}]]$  at  $\lambda 57 \mu\text{m}$  and its error.
- `OIII_88m` and `eOIII_88m`:  $[[\text{O III}]]$  at  $\lambda 88 \mu\text{m}$  and its error.
- `NII_122m` and `eNII_122m`:  $[[\text{N II}]]$  at  $\lambda 122 \mu\text{m}$  and its error.
- `NII_205m` and `eNII_205m`:  $[[\text{N II}]]$  at  $\lambda 205 \mu\text{m}$  and its error.

The program will then prompt for the model grid selection. For version 3.3, models from POPSTAR, BPASS, double composite AGN with  $\alpha_{ox}$  from -0.8 to -2.0, Rauch post-AGB for  $T_* = 50, 100$  or  $150 \text{ kK}$ , and ADAF models are available. In the case of these latter, the code will also ask for the assumed stopping criterion for the models (2% or 98% of free electrons). Additionally, as in other versions, users can include custom grids, by specifying them in the `Libraries_ir` folder. the code will prompt about model interpolation to enhance grid resolution by a factor of 10, though this may slow down calculations. Finally, the code will prompt for constraints when certain emission lines are missing. Again, the code can provide templates calculated for star-forming galaxies by Pérez-Montero, 2014, for EELGs by Pérez-Montero et al., 2021, or assuming low- or high-excitation AGNs. Additional constraint files can also be added by users, and the initial level of interactivity can be adjusted by setting the variable `interactive` in the script.

#### 4.2. Output File and Analysis of the Results

As with previous versions of HCM, after selecting the model grid, the code will display results for O/H, N/O and  $\log U$ . along with the ratio of completeness of the task. It will also create another `ascii` file whose first column is the identification of each row followed by the emission lines with their errors given as inputs. The next column indicates if the complete grid is used (index 1). However, as no auroral lines are available in this spectral range, this is not used in any case. The index 2 corresponds to the grid constrained following an empirical relation between O/H and  $\log U$ , as shown in Figure 1. Finally, index 3 is used when no previous estimation of N/O, using the N3O3 parameter can be done, and an empirical relation between O/H and N/O is assumed as shown in right panel of Figure 1. Alternatively, these relations can be changed to EELGs or others defined by the user, as commented above.

Table 3 lists mean offsets and standard deviations of the residuals for the resulting O/H, N/O, and  $\log U$  derived from the code from different sets of input emission lines, as compared with the values used as input in each grid of the model.

### 5. HCM for the Calculation of $T_{\text{eff}}$

The HII-CHI-MISTRY-TEFF (HCM-TEFF hereafter) differs from previous versions of HCM in its focus on calculating the effective temperature of the ionizing source ( $T_{\text{eff}}$ ) and/or the fraction of absorbed ionizing photons,  $f_{\text{abs}}$ , assuming a matter-bounded geometry. This tool uses the relationship between the 'softness parameter' and the hardness of the ionizing radiation (see Vilchez and Pagel, 1988, Pérez-Montero and Vilchez, 2009), as detailed in Pérez-Montero et al., 2019b and Pérez-Montero et al., 2020. A discussion on the impact of diffuse ionized gas (DIG) on results, particularly when low-excitation

TABLE 3. Mean offsets and standard deviation of the residuals of the resulting properties derived by HCm-IR when the models from POPSTAR emission lines are used as input as a function of the used emission-line ratios.

Used lines	Grid	O/H			N/O			$\log U$	
		Mean	$\Delta$	$\sigma$ res.	Mean	$\Delta$	$\sigma$ res.	Mean	$\Delta$
All lines	2	+0.04	0.09	-0.00	0.01	+0.02	0.07		
$[[\text{Ne II}]], [[\text{Ne III}]], [[\text{S II}]], [[\text{S IV}]], [[\text{O III}]], [[\text{N II}]], [[\text{N III}]]$	2	+0.03	0.13	+0.00	0.01	+0.02	0.07		
$[[\text{O III}]], [[\text{N II}]], [[\text{N III}]]$	2	-0.01	0.18	+0.00	0.01	+0.02	0.15		
$HI, [[\text{S III}]], [[\text{S IV}]], [[\text{Ne II}]], [[\text{Ne III}]]$	3	+0.04	0.04	-	-	-0.01	0.01		
$[[\text{S III}]], [[\text{S IV}]], [[\text{Ne II}]], [[\text{Ne III}]]$	3	+0.04	0.04	-	-	+0.00	0.01		
$[[\text{S III}]], [[\text{S IV}]]$	3	+0.02	0.02	-	-	+0.00	0.01		
$[[\text{Ne II}]], [[\text{Ne III}]]$	3	+0.03	0.09	-	-	+0.00	0.06		
$[[\text{N II}]], [[\text{N III}]]$	3	+0.00	0.11	-	-	+0.00	0.06		

lines like  $[[\text{S II}]]$  are used, is presented in Pérez-Montero et al., 2023b. An independent version for the infrared, valid for the derivation of  $T_{\text{eff}}$  and  $U$  in star-forming regions, is described in Pérez-Montero et al., 2024. This same version permits the derivation of the shape of the incident field of radiation in the NLR of AGN, as described in Pérez-Montero et al., 2025.

### 5.1. Running the Program and Preparing the Input File

Running the program is equivalent to the previous versions of HCm, through the terminal prompt:

```
> HCm-Teff_v5.2.py input.txt 50
```

The input file, required in `ascii` format, begins with the object ID and labels of the emission lines, including, if available, the oxygen abundance, which helps reduce uncertainty. The code accepts the following labels:

- ID: Identification name for each row.
- `12logOH` and `e12logOH`: Total oxygen abundance  $12+\log(\text{O/H})$  and its error. This may be calculated using HCm in a previous iteration.
- `CIV_1549` and `eCIV_1549`:  $[\text{C IV}] \lambda 1549 \text{ \AA}$  and its error.
- `CIII_1909` and `eCIII_1909`:  $[\text{C III}] \lambda 1909 \text{ \AA}$  and its error. This should include all lines taking part in the multiplet around this wavelength.
- `OII_3727` and `eOII_3727`:  $[\text{O II}] \lambda 3727 \text{ \AA}$  and its error. As in the case of HCm, this represents the addition of the two lines of the doublet of  $[\text{O II}]$  if there is good spectral resolution.
- `OIII_4959` and `eOIII_4959`:  $[\text{O III}] \lambda 4959 \text{ \AA}$  and its error.
- `OIII_5007` and `eOIII_5007`:  $[\text{O III}] \lambda 5007 \text{ \AA}$  and its error. As in the case of HCm, if only one of the two lines of the  $[\text{O III}]$  doublet (4959, 5007) is introduced, the code assumes its addition taking into account its theoretical relation.
- `HeI_4471` and `eHeI_4471`:  $[\text{He I}] \lambda 4471 \text{ \AA}$  and its error.
- `HeI_5876` and `eHeI_5876`:  $[\text{He I}] \lambda 5876 \text{ \AA}$  and its error.
- `HeI_6678` and `eHeI_6678`:  $[\text{He I}] \lambda 6678 \text{ \AA}$  and its error.
- `HeII_4686` and `eHeII_4686`:  $[\text{He II}] \lambda 4686 \text{ \AA}$  and its error.
- `SII_6716` and `eSII_6716`:  $[\text{S II}] \lambda 6716 \text{ \AA}$  and its error.

- **SII\_6731** and **eSII\_6731**:  $[[\text{S II}]] \lambda 6731 \text{ \AA}$  and its error. It is also possible to use the addition of the two  $[[\text{S II}]]$  lines using **SII\_6725** and **eSII\_6725**.
- **SIII\_9069** and **eSIII\_9069**:  $[[\text{S III}]] \lambda 9069 \text{ \AA}$  and its error.
- **SIII\_9532** and **eSIII\_9532**:  $[[\text{S III}]] \lambda 9532 \text{ \AA}$  and its error. As in the case of  $[[\text{O III}]]$ , if one of the two  $[[\text{S III}]]$  nebular lines is not introduced, the code assumes its addition taking the theoretical expected ratio.
  - **ArIV\_4740** and **eArIV\_4740**:  $[\text{ArIV}] \lambda 4740 \text{ \AA}$  and its error.
  - **ArIII\_7135** and **eArIII\_7135**:  $[\text{ArIII}] \lambda 7135 \text{ \AA}$  and its error.
  - **NII\_6584** and **eNII\_6584**:  $[[\text{N II}]] \lambda 6584 \text{ \AA}$  and its error.

The next rows in the input file should correspond to the different values for which the code should derive  $T_{\text{eff}}$  and  $\log U$ . The provided emission line fluxes must be reddening corrected, but it is not necessary that they are relative to  $\text{H}\beta$ , as only ratios of lines of different ionization are used. Use 0 when no value is found for a line or error.

### 5.2. The Grids of Models

If the input file is correct and the variable `interactive` is set to `True` in the code, it will prompt the parameters we want to calculate:

- (1) **Effective temperature and ionization parameter**
- (2) **Photon absorption fraction and ionization parameter**

**Choose derived parameters:**

Depending on these parameters, the code will prompt for different grids of models and geometries. In the case of  $T_{\text{eff}}$ :

- (1) **WM-Basic (30-60 kK)**
- (2) **WM-Basic (30-60 kK) and Rauch (80-120 kK) stellar atmospheres**
- (3) **Black body (30-100 kK)**

**Choose models:**

Option 1 calculates  $T_{\text{eff}}$  and  $\log U$  using WM-BASIC single star atmospheres from [Paul-drach et al., 2001](#) from 30 to 60 kK, while option 2 extends the range up to 120 kK using post-AGB stellar atmospheres from Rauch. Option 3 uses black-body spectral energy distributions in the range 30-100 kK.

After selecting the model type, the program will prompt for the assumed geometry, given its impact on model positions within softness diagrams:

- (1) **Plane-parallel geometry**
- (2) **Spherical geometry**

**Choose geometry of the models:**

If we want to calculate the photon absorption factor,  $f_{\text{abs}}$ , defined as the ratio of ionizing hydrogen photons that do not escape from the nebula, along with  $\log U$ , this is calculated using cluster model atmospheres from BPASS v.2.1 from [Eldridge et al., 2017](#), assuming an instantaneous burst at 4 Myr, with binaries and an IMF with slope  $X = -1.35$  and an upper mass limit of  $300 \text{ M}_{\odot}$ . In this case, we can choose among models with stellar metallicity identical to that of the gas in each model, or either assume nearly metal-free stars ( $Z = 10^{-5}$ ). The code also supplies models with black-body with  $T_{\text{eff}} =$

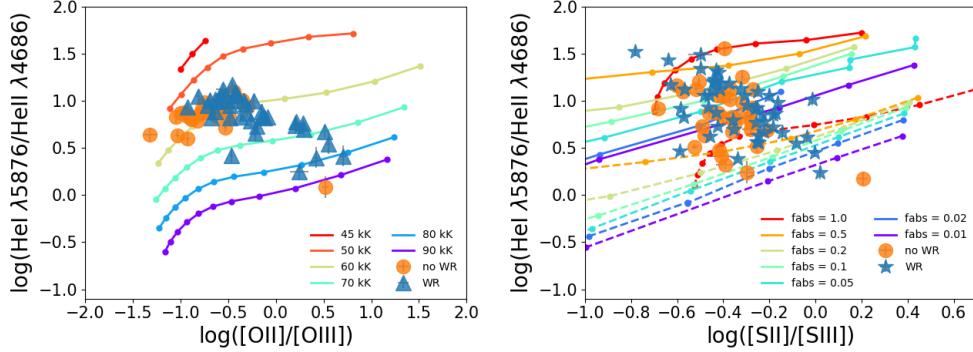


FIGURE 2. Examples of two softness diagrams for the sample of He II-emitters presented in Pérez-Montero et al., 2020. At left  $[\text{O II}]/[\text{O III}]$  vs  $\text{HeI}/\text{He II}$  with models at  $12 + \log(\text{O/H})$  and different values for  $T_{\text{eff}}$  and, at right, using  $[\text{S II}]/[\text{S III}]$  and BPASS models assuming different values for  $f_{\text{abs}}$ . In all sequences values in the lower left part correspond to higher values for  $\log U$ .

$10^5$  K, equivalent for metal-free stars.

- (1) BPASS cluster atmospheres, age = 4 Myr, Mup = 300, x = 1.35, w/binaries,  $Z_* = Z_g$
- (2) BPASS cluster atmospheres, age = 4 Myr, Mup = 300, x = 1.35, w/binaries,  $Z_* = 1e-5$
- (3) Black-body,  $T_* = 1e5$  K

In this case, only spherical geometry models are considered.

In Figure 2, two examples of the behavior of the grids of models in one of the softness diagrams are shown, as in Pérez-Montero et al., 2020, illustrating how the grids of models cover the space of the emission line ratios.

Finally, the code will prompt for the possibility of using interpolation in the final chosen grid:

Choose

### 5.3. Output file and analysis of the results

As in other versions of HCm, the code will show on the screen the results for each row of the input file along with the ratio of completeness of the task. At the end it will create an output file called as the name of the input file adding `hcm-teff--output.dat`. This file contains all the information for the assumed grids and the results for each one of the rows.

The first column lists the row identification, followed by input fluxes for each emission line, then the assumed O/H value (set to 0 if absent), and finally, the derived  $T_{\text{eff}}$  (or  $f_{\text{abs}}$  if Grid 3 is selected) and  $\log U$  with associated errors. A value of 0 indicates that no solution was found due to insufficient line data, while -99999 signals an error with the specific input row.

Table 4 lists the mean offsets and the standard deviation of the residuals for the obtained final results as a function of the input emission lines when compared with the input information from the code.

TABLE 4. Mean offsets and standard deviation of the residuals of the resulting properties derived by HCM-TEFF when the model emission lines are used as input as a function of the used emission-line ratios.

Used line ratios	$T_{\text{eff}}$ (kK) <sup>†</sup>			$\log U^a$			$\log f_{\text{abs}}^{\ddagger}$			$\log U^b$		
	Mean	$\Delta$	$\sigma$ res.	Mean	$\Delta$	$\sigma$ res.	Mean	$\Delta$	$\sigma$ res.	Mean	$\Delta$	$\sigma$ res.
$[[\text{O II}]]/[[\text{O III}]]$ , $[[\text{S II}]]/[[\text{S III}]]$ , $\text{He I}/\text{He II}$	+0.8	2.1	+0.01	0.10	+0.11	0.34	-0.10	0.53				
$[[\text{O II}]]/[[\text{O III}]]$ , $\text{He I}/\text{He II}$	+0.8	2.1	+0.01	0.10	-0.12	0.27	+0.14	0.47				
$[[\text{S II}]]/[[\text{S III}]]$ , $\text{He I}/\text{He II}$	+0.6	2.0	+0.03	0.13	-0.19	0.35	+0.10	0.62				
$[[\text{S II}]]/[[\text{O III}]]$ , $\text{He I}/\text{He II}$	-0.3	5.1	+0.04	0.51	-0.07	0.27	-0.10	0.59				
$\text{He I}/\text{He II}$	-0.5	6.8	+0.04	0.72	-0.12	0.35	-0.25	0.63				

#### 5.4. The HCM-TEFF-IR version

A similar code is available for its use based only on mid-IR emission-lines. Although it is planned to merge it with the optical version, so far it works in an independent module. It is described and applied in Pérez-Montero et al., 2024 to calculate  $T_{\text{eff}}$  for star-forming regions and in Pérez-Montero et al., 2025 to calculate  $\alpha_{OX}$  in AGN, taking advantage of the high-excitation lines that can be measured in this regime.

It basically works in the same way as it is described above, using exactly the same grids of models with the exception of those calculated for the derivation of the photon absorption factor. New models have been also calculated for AGN in order to cover a wide range for  $\alpha_{OX}$  in the range [-2.0, -1.0] and  $\log U$  in the range [-4.0, -1.0] for oxygen abundances from 12+log(O/H) from 8.1 to 9.0. The user can select for this modality between models with or without dust grains mixed in the gas and test different stopping criteria in the models (2%, 98%, and 99.9% of free electrons) what gives the possibility of exploring the weight in the results of the low-excitation zone in the integrated spectrum.

The instructions and steps are the same, but using in the input the following emission-lines, in addition to the ID and metallicity entries:

- ArIII\_7m and ArIII\_7m for  $[[\text{Ar III}]]$  at  $\lambda$  6.98  $\mu\text{m}$  and its error.
- ArIV\_8m and ArIV\_8m for  $[[\text{Ar IV}]]$  at  $\lambda$  7.90  $\mu\text{m}$  and its error.
- ArIII\_9m and ArIII\_9m for  $[[\text{Ar III}]]$  at  $\lambda$  8.99  $\mu\text{m}$  and its error.
- SIV\_10m and eSIV\_10m for  $[[\text{S IV}]]$  at  $\lambda$  at 10.5  $\mu\text{m}$  and its error
- NeII\_12m and eNeII\_12m for  $[[\text{Ne II}]]$   $\lambda$  12.8  $\mu\text{m}$  and its error.
- ArV\_13m and ArV\_13m for  $[[\text{Ar V}]]$  at  $\lambda$  13.1  $\mu\text{m}$  and its error.
- NeV\_14m and eNeV\_14m for  $[\text{Ne V}]$   $\lambda$  14.9  $\mu\text{m}$  and its error.
- NeIII\_15m and eNeIII\_15m for  $[[\text{Ne III}]]$   $\lambda$  15.5  $\mu\text{m}$  and its error.
- SIII\_18m and eSIII\_18m for  $[[\text{S III}]]$   $\lambda$  18.8  $\mu\text{m}$  and its error.
- NeV\_24m and eNeV\_24m for  $[\text{Ne V}]$   $\lambda$  24.3  $\mu\text{m}$  and its error.
- OIV\_25m and eOIV\_25m for  $[\text{O IV}]$   $\lambda$  25.9  $\mu\text{m}$  and its error.
- SIII\_33m and eSIII\_33m for  $[[\text{S III}]]$   $\lambda$  at 33.7  $\mu\text{m}$  and its error.
- OIII\_52m and OIII\_52m for  $[[\text{O III}]]$   $\lambda$  52  $\mu\text{m}$  and its error.
- NII\_57m and eNII\_57m for  $[[\text{N II}]]$   $\lambda$  57  $\mu\text{m}$  and its error
- OIII\_88m and eOIII\_88m for  $[[\text{O III}]]$   $\lambda$  88  $\mu\text{m}$  and its error
- NII\_122m and eNII\_122m for  $[[\text{N II}]]$   $\lambda$  122  $\mu\text{m}$  and its error.
- NII\_205m and eNII\_205m for  $[[\text{N II}]]$   $\lambda$  205  $\mu\text{m}$  and its error.

In this case, the code uses as observables ratios of low- to -high-ionization IR emission-lines that can also be used to construct the corresponding IR softness diagram (e.g.  $[[\text{Ne II}]]/[[\text{Ne III}]]$ ,  $[[\text{S III}]]/[[\text{S IV}]]$ ). As compared to the optical version, this version is less affected by DIG contribution as it involves higher excitation emission-lines.

The very high excitation lines, such as  $[[\text{O IV}]]$  or  $[[\text{Ne V}]]$  can also be used to discriminate very hard incident SEDs. In the case of AGN the line ratios can only be used to provide an estimate of the hardness when they involve very high-excitation lines (e.g.  $([[\text{Ne II}]]+[[\text{Ne III}]])/[[\text{Ne V}]]$  or  $[[\text{O III}]]/[[\text{O IV}]]$ , or  $([[\text{S III}]]+[[\text{S IV}]])/[[\text{O IV}]]$ ).

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